



Bose-Einstein Condensation of Atoms in a Uniform Potential

Alexander L. Gaunt, Tobias F. Schmidutz, Igor Gotlibovych, Robert P. Smith, and Zoran Hadzibabic
Cavendish Laboratory, University of Cambridge, J. J. Thomson Avenue, Cambridge CB3 0HE, United Kingdom
 (Received 14 March 2013; published 16 May 2013)

We have observed the Bose-Einstein condensation of an atomic gas in the (quasi)uniform three-dimensional potential of an optical box trap. Condensation is seen in the bimodal momentum distribution and the anisotropic time-of-flight expansion of the condensate. The critical temperature agrees with the theoretical prediction for a uniform Bose gas. The momentum distribution of a noncondensed quantum-degenerate gas is also clearly distinct from the conventional case of a harmonically trapped sample and close to the expected distribution in a uniform system. We confirm the coherence of our condensate in a matter-wave interference experiment. Our experiments open many new possibilities for fundamental studies of many-body physics.

DOI: 10.1103/PhysRevLett.110.200406

PACS numbers: 03.75.Hh, 67.85.-d

Ultracold Bose and Fermi atomic gases are widely used as test beds of fundamental many-body physics [1]. Experimental tools such as Feshbach interaction resonances [2], optical lattices [3], and synthetic gauge fields [4] offer great flexibility for studies of outstanding problems arising in many areas, most commonly in condensed-matter physics. However, an important difference between “conventional” many-body systems and ultracold gases is that the former are usually spatially uniform whereas the latter are traditionally produced in harmonic traps with no translational symmetries.

Various methods have been developed to overcome this problem and extract uniform-system properties from a harmonically trapped sample [5–13], relying on the local density approximation [5–11] or selective probing of a small central portion of the cloud [11–13]. Sometimes harmonic trapping can even be advantageous, allowing simultaneous mapping of uniform-system properties at different (local) particle densities. On the other hand, in many important situations local approaches are inherently limiting, for example, for studies of critical behavior with diverging correlation lengths. The possibility to directly study a spatially uniform quantum-degenerate gas has thus remained an important experimental challenge. So far, atomic Bose-Einstein condensates (BECs) have been loaded into elongated [14] or toroidal [15] traps that are uniform along only one direction while still harmonic along the other two directions.

Here, we demonstrate the Bose-Einstein condensation of an atomic gas in a three-dimensional (3D) (quasi)uniform potential. We load an optical box trap depicted in Fig. 1(a) with ^{87}Rb atoms pre-cooled in a harmonic trap and achieve condensation by evaporative cooling in the box potential. Below a critical temperature $T_c \approx 90$ nK, condensation is seen in the emergence of a bimodal momentum distribution and the anisotropic time-of-flight (TOF) expansion of the BEC. We characterize the flatness of our box potential and show that both the momentum distribution of the

non-condensed component and the thermodynamics of condensation are close to the theoretical expectations for a uniform system, while being clearly distinct from the conventional case of a harmonically trapped gas. We also

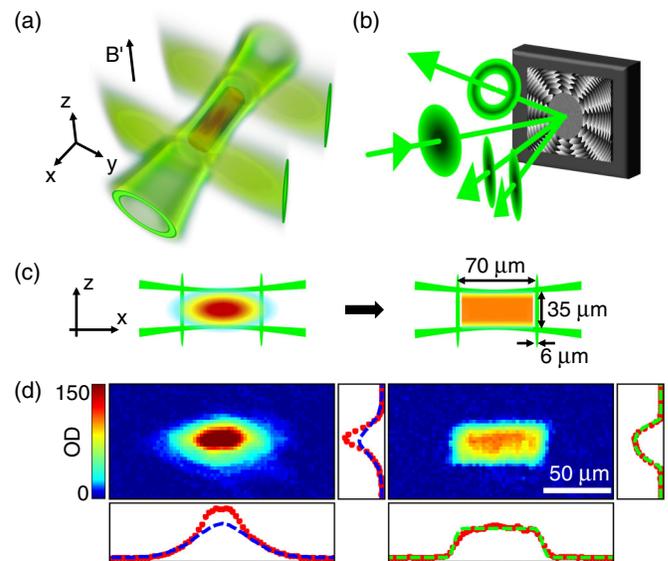


FIG. 1 (color online). Preparing a quasiuniform Bose gas. (a) The optical-box trap is formed by one hollow tube beam and two sheet beams creating a repulsive potential for the atoms. The atomic cloud is confined to the dark (red) cylindrical region. Gravitational force is canceled by a magnetic field gradient B' . (b) The three trapping beams are created by reflecting a single Gaussian beam off a phase-imprinting spatial light modulator. (c) The atoms are loaded into the box trap after precooling in a harmonic trap. (d) *In situ* images of the cloud just before (left) and after (right) loading into the box and corresponding line-density profiles along x (bottom plots) and z (side plots) directions. OD stands for optical density; the line densities along x (z) are obtained by integrating the images along z (x). The blue dashed lines in the left panel are fits to the thermal component of the harmonically trapped gas. The green dashed lines in the right panel are fits based on the expected profiles for a uniform-density gas.

demonstrate a simple experimental configuration for trapped-atom interferometry and use it to confirm the coherence of our quasi-uniform BEC.

Our setup for producing ^{87}Rb condensates in a harmonic potential is described elsewhere [16]; we create BECs in the $|F, m_F\rangle = |2, 2\rangle$ hyperfine ground state using a hybrid magnetic-optical trap [17]. The dark optical trap [18,19] that is central to this work is formed by three 532 nm laser beams—a “tube” beam propagating along the x axis and two “sheet” beams propagating along the y axis. The green laser beams create a repulsive potential for the atoms and confine them to the cylindrical dark region depicted in red in Fig. 1(a). To create a uniform potential, we additionally cancel the gravitational force on the atoms at a 10^{-4} level, using a magnetic field gradient [16].

As outlined in Fig. 1(b), all three trapping beams are created by reflecting a single Gaussian beam off a phase-imprinting spatial light modulator with three superposed phase patterns [20]. The tube beam is an optical vortex created by imprinting a 24π phase winding on the incoming beam [21], the sheet beams are created using cylindrical-lens phase patterns, and the three outgoing beams are deflected in different directions using phase gradients. With a total laser power of $P_0 \approx 700$ mW we achieve a trap depth of $V_0 \approx k_B \times 2 \mu\text{K}$.

We evaporatively cool the gas in the harmonic trap down to $T \approx 120$ nK, when the cloud size is similar to the size of our optical box [see Fig. 1(c)] and $k_B T \ll V_0$. At this point the gas is partially condensed, but the BEC is lost during the transfer into the box trap, which is not perfectly adiabatic. Over 1 s, we turn on the green light and then turn off the harmonic trapping, capturing $>80\%$ of the atoms.

In Fig. 1(d) we show in-trap absorption images of the cloud just before and just after the transfer into the box trap. The images are taken along the y direction, using high-intensity imaging [8,9,22] with a saturation parameter $I/I_{\text{sat}} \approx 150$. For each image, we show the line-density profiles along x and z , obtained by integrating the image along one direction. If a cylindrical box of length L and radius R is filled perfectly uniformly, the density distribution along x is simply a top-hat function of width L . Along z , the line-of-sight integration results in “circular” column- and line-density profiles, $\propto \sqrt{1 - (z/R)^2}$. In the experimental images, the edges of the cloud are rounded off for two reasons, both related to the diffraction limit of our optical setup. First, the $1/e^2$ waist of the 532 nm trapping beams is diffraction limited to $\approx 3 \mu\text{m}$, which leads to some rounding off of the potential bottom near the edges of the box. Second, our imaging resolution is diffraction limited to $\approx 5 \mu\text{m}$, making the cloud edges appear more smeared out than they actually are. The green dashed lines in the right panel of Fig. 1(d) are fits to the data based on a perfectly uniform distribution convolved *only* with the imaging point-spread function. The fits describe the data well and give $L = 63 \pm 2 \mu\text{m}$ and $R = 15 \pm 1 \mu\text{m}$. These

values are consistent with the calculated separation of the green walls, reduced by the diffraction-limited wall thickness.

After the transfer into the box trap, the cloud contains $N \approx 6 \times 10^5$ atoms at $T \approx 130$ nK. From this point, we cool the gas to below T_c by forced evaporative cooling in the box trap. We lower the trapping power P in an exponential ramp with a 0.5 s time constant, thus proportionally reducing the power in all three trapping beams. Initially, the trap depth is much larger than $k_B T$, so significant cooling occurs only for $P \lesssim 0.5P_0$. At the end of the evaporation, we always raise (over 0.5 s) the trapping power back to P_0 so that the cloud cooled to different temperatures is always confined in the same potential.

Figure 2 qualitatively illustrates the effects of evaporation and condensation in the box trap. We show images of the cloud both *in situ* and after 50 ms of TOF expansion from the trap. Whereas in a harmonic trap cooling results in simultaneous real-space and momentum-space condensation, here it has no dramatic effects on the in-trap atomic distribution. The density is gradually reduced by evaporation, but the shape of the cloud does not reveal condensation. On the other hand, in momentum space (i.e., in TOF) the effects of cooling are obvious and the signatures of condensation are qualitatively the same as for a harmonically trapped gas—the momentum distribution becomes bimodal and the BEC expands anisotropically, with its aspect ratio inverting in TOF.

We now turn to a quantitative analysis of our degenerate quasiuniform Bose gas. We assess the flatness of our trapping potential and contrast the thermodynamics of condensation in our system with the case of a harmonically trapped gas.

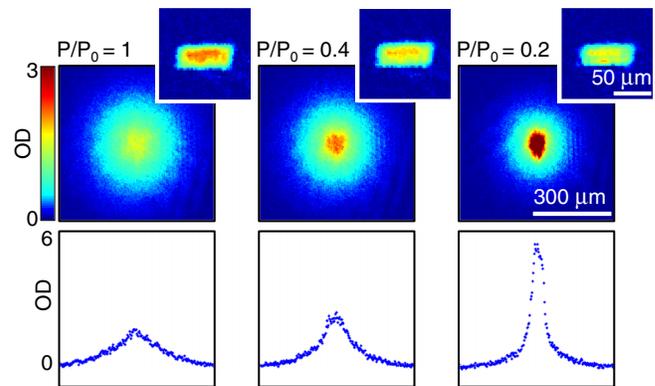


FIG. 2 (color online). Evaporation and Bose-Einstein condensation in the optical-box trap. Cooling is achieved by lowering the trapping laser power P . We show absorption images taken after 50 ms of TOF and *in situ* [insets, with same color scale as in Fig. 1(d)]. The bottom panels show cuts through the momentum distributions recorded in TOF. In contrast to the case of a harmonic trap, no dramatic effects of cooling are observed *in situ*. However, BEC is clearly seen in the bimodality of the momentum distribution and the anisotropic expansion.

We first analyze the momentum distribution in a noncondensed gas, at $T \approx 110$ nK. In Fig. 3(a), we plot the line density measured in TOF, $\tilde{n}(z)$, obtained by integrating an image along x . At this point, the gas is sufficiently degenerate to show the effects of Bose statistics. In a trapped degenerate gas, the spatial and momentum degrees of freedom are in general not separable, and consequently $\tilde{n}(z)$ also contains information about the functional form of the trap.

For simplicity, let us consider an isotropic 3D trapping potential of the form $V(r) \propto r^n$; for a fixed n value, potential anisotropy does not affect the scalings discussed below. The line-density distribution after long TOF then has the form

$$\tilde{n}(z) = \sqrt{\frac{\beta m}{2t^2}} \left(\frac{T}{\theta_\alpha}\right)^{\alpha+1/2} g_\alpha(e^{\beta(\mu - \varepsilon(z))}), \quad (1)$$

where g_α is the polylog function of order $\alpha = 1 + 3/n$. Here, $\beta = 1/(k_B T)$, m is the atom mass, t is the expansion time, μ is the chemical potential, $\varepsilon(z) = mz^2/(2t^2)$, and the constant θ_α absorbs various factors such as the imaging magnification and cross section. Small corrections due to the initial cloud size are accounted for by convolving g_α with the in-trap density distribution.

For a harmonic trap $n = 2$ and $\alpha = 5/2$, whereas for a uniform gas $n \rightarrow \infty$ and $\alpha = 1$. The measured \tilde{n} is fitted very well by the g_1 function, with the reduced $\chi^2 \approx 1$. The $g_{5/2}$ fit is comparatively poor: $\chi^2 \approx 7$ and the systematic patterns in the fitting residuals $\Delta\tilde{n} = \tilde{n} - \tilde{n}_{\text{fit}}$ clearly show that this is fundamentally a wrong functional form. Qualitatively, the measured momentum distribution is more “peaky” than that of a harmonically trapped degenerate gas.

As shown in the inset of Fig. 3(a), the measured \tilde{n} can be fitted well using α in the range of 1–1.7, whereas higher values can be clearly excluded. Note, however, that in these fits we used θ_α , μ , and T as free parameters. Allowing θ_α to vary gives the fitting function an unphysical freedom and overestimates the range of suitable α values. Crucially, *only* for the correct α is the best-fit value of θ_α a temperature-independent constant. We use this fact to accurately determine the leading-order correction to the flatness of our trapping potential.

In Fig. 3(b), we analyze the T dependence of the fitted θ_α values. (As above, T is varied by forced evaporation in the box trap, and at the end of cooling the trap is always restored to its full depth.) Here, we exclude from the fits the central 180 μm wide region of the cloud, which is larger than the largest observed BECs. The drift of the fitted θ_α with T is described well by a constant slope $S_\alpha = d[\log(\theta_\alpha)]/d[\log(T)]$, and we see that S_α varies monotonically with α . Physically, if the α used for fitting is too high, as the gas cools the actual momentum distribution narrows faster than g_α . The fit then increasingly underestimates T and compensates by decreasing θ_α . Conversely, if α is too low, the fitted θ_α increases as the gas is cooled. Note that for each α the fitted θ_α is plotted against T extracted from the same fit, and we see that the fitted temperatures also show the expected systematic drifts.

From the condition $S_\alpha = 0$, we get $\alpha = 1.23 \pm 0.03$ and conclude that the leading-order correction to the flatness of our box potential is $\propto r^{13 \pm 2}$. While we can distinguish an r^{13} potential from a perfectly flat one, we expect this distinction to be irrelevant for most many-body studies.

We now fix $\alpha = 1.23$ and study the evolution of the BEC atom number N_0 with T . In contrast to the harmonic-trap case, where the Thomas-Fermi shape of the BEC

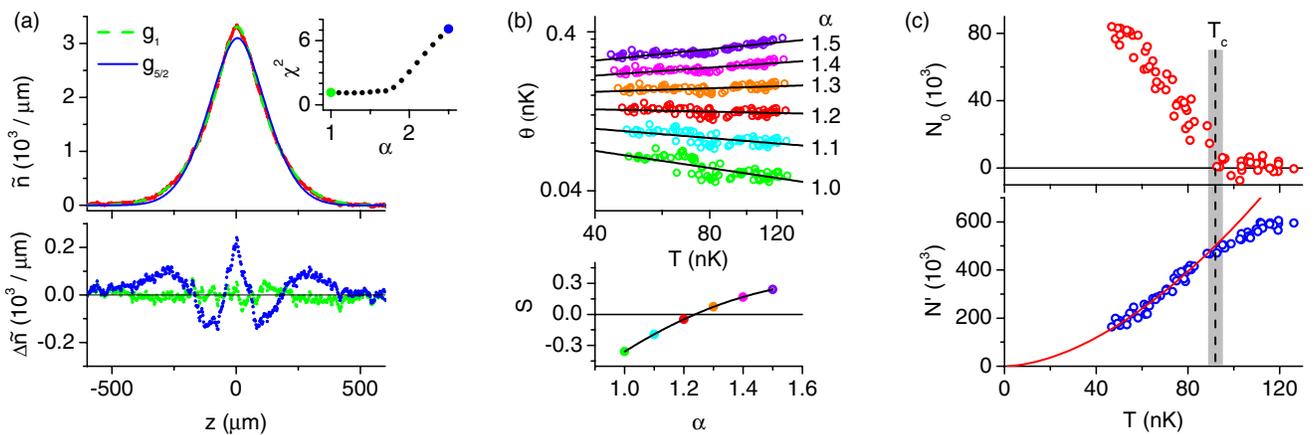


FIG. 3 (color online). Momentum and particle-number distribution in a quasiuniform degenerate Bose gas. (a) TOF profile $\tilde{n}(z)$ of a noncondensed cloud (red) and g_α fits to the data, with $\alpha = 1$ (green) and $5/2$ (blue). We also show the fitting residuals $\Delta\tilde{n}$. Inset: fitting χ^2 versus α . (b) Flatness of the box potential: the condition $S = d[\log(\theta)]/d[\log(T)] = 0$ (see text) gives $\alpha = 1.23 \pm 0.03$, corresponding to an effective $r^{13 \pm 2}$ potential. (c) The number of condensed (N_0) and thermal (N') atoms versus T . The critical temperature $T_c = 92 \pm 3$ nK is in agreement with the prediction for a uniform Bose gas. Below T_c , a power-law fit (solid red line) gives $N' \propto T^{1.73 \pm 0.06}$.

remains parabolic in TOF, the shape of the BEC expanding from our uniform cylindrical trap does not maintain a simple functional form. Hence, rather than fitting the BEC profile to any specific shape, we fit only the thermal component and simply count the excess atoms not accounted for by the thermal fit.

In Fig. 3(c) we plot N_0 versus T and determine $T_c = 92 \pm 3$ nK. From the measured atom number and box volume, we get a consistent theoretical value for a uniform Bose gas, $T_c^0 = 98 \pm 10$ nK. This calculation includes small finite-size and interaction shifts of T_c^0 [23–25], of about 6% and 1%, respectively; the error includes the 10% uncertainty in our absolute atom-number calibration, obtained by measuring T_c in a harmonic trap [26].

In Fig. 3(c) we also plot the thermal atom number N' versus T . Below T_c , we fit the variation of N' with a power law $N' \propto T^\gamma$ and get $\gamma = 1.73 \pm 0.06$. This agrees with the expected $\gamma = \alpha + 1/2$. For a saturated thermal component in a perfectly uniform system $N' \propto T^{3/2}$, whereas in a harmonic trap $N' \propto T^3$. The thermodynamics of our gas are therefore very close to the textbook case of a uniform system and very different from the case of a harmonically trapped sample.

Below 40 nK our temperature fits are not accurate, but we can cool the gas further and produce a quasipure BEC with $>10^5$ atoms. The $1/e$ lifetime of the BEC is 10 s, limited by the background pressure in our vacuum chamber [16].

Finally, we verify the coherence of our uniform BEC in an interference experiment, using a two-trap configuration shown in Fig. 4(a). We add a third sheet laser beam to simultaneously condense atoms in a small “satellite” trap, whereas the main box trap is essentially the same as in the rest of the Letter. Here, we cool the gas to the regime where the BEC is quasi-pure ($T < 30$ nK).

In Fig. 4(b) we show the interference fringes emerging as the two clouds overlap in TOF. We observe interference contrasts of up to $\approx 50\%$, limited by our imaging resolution and the fact that the two interfering clouds do not have the same atom number. Note that the interference pattern is different from the parallel fringes of constant spacing commonly observed in a symmetric expansion of two BECs [27]. Qualitatively, along the x axis the main cloud expands very slowly and is still in the near-field regime, while the small BEC expands rapidly and is in the far-field regime. This leads to a quadratic variation of the relative phase along x and the reduction of the fringe period from right to left in Fig. 4(b). We reproduce a similar pattern in a numerical simulation of the Gross-Pitaevskii equation, shown in Fig. 4(c).

In conclusion, we have observed and characterized the Bose-Einstein condensation of an atomic gas in an essentially uniform potential. We expect our experiments to open up many new research possibilities. The box trap eliminates the need to rely on the local density approximation, which

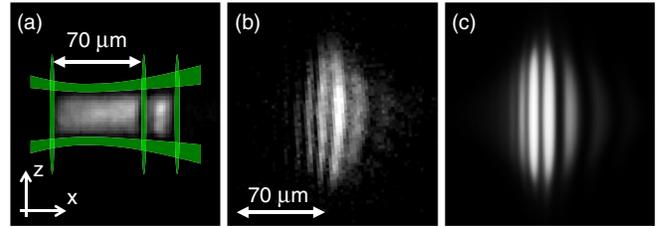


FIG. 4 (color online). Matter-wave interference with a quasiuniform BEC. (a) *In situ* image of a two-trap configuration, together with a sketch of the green-light arrangement used to create it. (b) Interference pattern observed after 100 ms of TOF expansion. (c) Numerical simulation of the expected interference pattern.

could be particularly important for studies of critical behavior with diverging correlations near phase transitions [28,29]. Moreover, for the same particle number, the quasiuniform BEC has a significantly lower density (in our case $\sim 2 \times 10^{12}$ cm $^{-3}$) than a harmonically trapped one. This reduces the importance of three-body recombination compared to that of two-body interactions, and near a Feshbach resonance it may facilitate equilibrium studies of a strongly interacting, possibly unitary Bose gas [30–32]. Dark optical traps loaded with low density BECs also hold promise for trapped-atom interferometry with very long decoherence times. Here, we chose an elongated box shape for qualitative detection of the BEC through anisotropic expansion, but with our light-shaping methods traps of almost arbitrary geometry could be created [33]. Our methods are also suitable for studies of degenerate Fermi gases and low-dimensional systems and are compatible with the implementation of 3D optical lattices.

This work was supported by EPSRC (Grant No. EP/G026823/1), the Royal Society, AFOSR, ARO, and DARPA OLE. We thank Laser Quantum for the loan of the 532 nm laser.

-
- [1] I. Bloch, J. Dalibard, and W. Zwerger, *Rev. Mod. Phys.* **80**, 885 (2008).
 - [2] C. Chin, R. Grimm, P. Julienne, and E. Tiesinga, *Rev. Mod. Phys.* **82**, 1225 (2010).
 - [3] O. Morsch and M. Oberthaler, *Rev. Mod. Phys.* **78**, 179 (2006).
 - [4] J. Dalibard, F. Gerbier, G. Juzeliūnas, and P. Öhberg, *Rev. Mod. Phys.* **83**, 1523 (2011).
 - [5] T.-L. Ho and Q. Zhou, *Nat. Phys.* **6**, 131 (2010).
 - [6] S. Nascimbène, N. Navon, K.J. Jiang, F. Chevy, and C. Salomon, *Nature (London)* **463**, 1057 (2010).
 - [7] N. Navon, S. Nascimbène, F. Chevy, and C. Salomon, *Science* **328**, 729 (2010).
 - [8] C.-L. Hung, X. Zhang, N. Gemelke, and C. Chin, *Nature (London)* **470**, 236 (2011).
 - [9] T. Yefsah, R. Desbuquois, L. Chomaz, K.J. Günter, and J. Dalibard, *Phys. Rev. Lett.* **107**, 130401 (2011).

- [10] M. J. H. Ku, A. T. Sommer, L. W. Cheuk, and M. W. Zwierlein, *Science* **335**, 563 (2012).
- [11] R. P. Smith, N. Tammuz, R. L. D. Campbell, M. Holzmann, and Z. Hadzibabic, *Phys. Rev. Lett.* **107**, 190403 (2011).
- [12] T. E. Drake, Y. Sagi, R. Paudel, J. T. Stewart, J. P. Gaebler, and D. S. Jin, *Phys. Rev. A* **86**, 031601 (2012).
- [13] Y. Sagi, T. E. Drake, R. Paudel, and D. S. Jin, *Phys. Rev. Lett.* **109**, 220402 (2012).
- [14] T. P. Meyrath, F. Schreck, J. L. Hanssen, C.-S. Chuu, and M. G. Raizen, *Phys. Rev. A* **71**, 041604 (2005).
- [15] S. Gupta, K. W. Murch, K. L. Moore, T. P. Purdy, and D. M. Stamper-Kurn, *Phys. Rev. Lett.* **95**, 143201 (2005).
- [16] I. Gotlibovych, T. F. Schmidutz, S. Moulder, R. L. D. Campbell, N. Tammuz, R. J. Fletcher, A. L. Gaunt, S. Beattie, R. P. Smith, and Z. Hadzibabic, [arXiv:1212.4108](https://arxiv.org/abs/1212.4108).
- [17] Y.-J. Lin, A. R. Perry, R. L. Compton, I. B. Spielman, and J. V. Porto, *Phys. Rev. A* **79**, 063631 (2009).
- [18] A. Kaplan, N. Friedman, and N. Davidson, *J. Opt. Soc. Am. B* **19**, 1233 (2002).
- [19] A. Jaouadi, N. Gaaloul, B. Viaris de Lesegno, M. Telmini, L. Pruvost, and E. Charron, *Phys. Rev. A* **82**, 023613 (2010).
- [20] J. Liesener, M. Reicherter, T. Haist, and H. J. Tiziani, *Opt. Commun.* **185**, 77 (2000).
- [21] Theoretically, at energies negligible compared to the trap depth this corresponds to a power-law potential $\propto r^{24}$ in the y - z plane. However, when a cloud explores a finite fraction of the trap depth, this reduces the effective power-law exponent seen in thermodynamics measurements.
- Generally, higher phase winding makes the potential flatter, but for a given focusing of the beam (limited by the optical access to the atoms) it also results in a larger box and lower T_c for the same atom number. We chose the 24π winding as an empirical compromise, giving a very flat potential while keeping T_c sufficiently high for accurate measurements.
- [22] G. Reinaudi, T. Lahaye, Z. Wang, and D. Guéry-Odelin, *Opt. Lett.* **32**, 3143 (2007).
- [23] S. Grossmann and M. Holthaus, *Z. Phys. B* **97**, 319 (1995).
- [24] P. Arnold and G. Moore, *Phys. Rev. Lett.* **87**, 120401 (2001).
- [25] V. A. Kashurnikov, N. V. Prokof'ev, and B. V. Svistunov, *Phys. Rev. Lett.* **87**, 120402 (2001).
- [26] R. P. Smith, R. L. D. Campbell, N. Tammuz, and Z. Hadzibabic, *Phys. Rev. Lett.* **106**, 250403 (2011).
- [27] M. R. Andrews, C. G. Townsend, H.-J. Miesner, D. S. Durfee, D. M. Kurn, and W. Ketterle, *Science* **275**, 637 (1997).
- [28] T. Donner, S. Ritter, T. Bourdel, A. Öttl, M. Köhl, and T. Esslinger, *Science* **315**, 1556 (2007).
- [29] P. Cladé, C. Ryu, A. Ramanathan, K. Helmerson, and W. D. Phillips, *Phys. Rev. Lett.* **102**, 170401 (2009).
- [30] N. Navon, S. Piatecki, K. Günter, B. Rem, T. C. Nguyen, F. Chevy, W. Krauth, and C. Salomon, *Phys. Rev. Lett.* **107**, 135301 (2011).
- [31] R. J. Wild, P. Makotyn, J. M. Pino, E. A. Cornell, and D. S. Jin, *Phys. Rev. Lett.* **108**, 145305 (2012).
- [32] W. Li and T.-L. Ho, *Phys. Rev. Lett.* **108**, 195301 (2012).
- [33] A. L. Gaunt and Z. Hadzibabic, *Sci. Rep.* **2**, 721 (2012).